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Propagation Failure In The Discrete Nagumo Equation



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Lattice Differential Equations

Lattice differential equations (LDEs) are ODEs indexed on a spatial lattice, e.g.

$$\frac{d}{dt}u_j(t) = \alpha(u_{j-1}(t) - 2u_j(t) + u_{j+1}(t)) + f(u_j(t)), \quad j \in \mathbb{Z}.$$



Picking $\alpha = h^{-2} \gg 1$, LDE can be seen as discretization with distance h of PDE

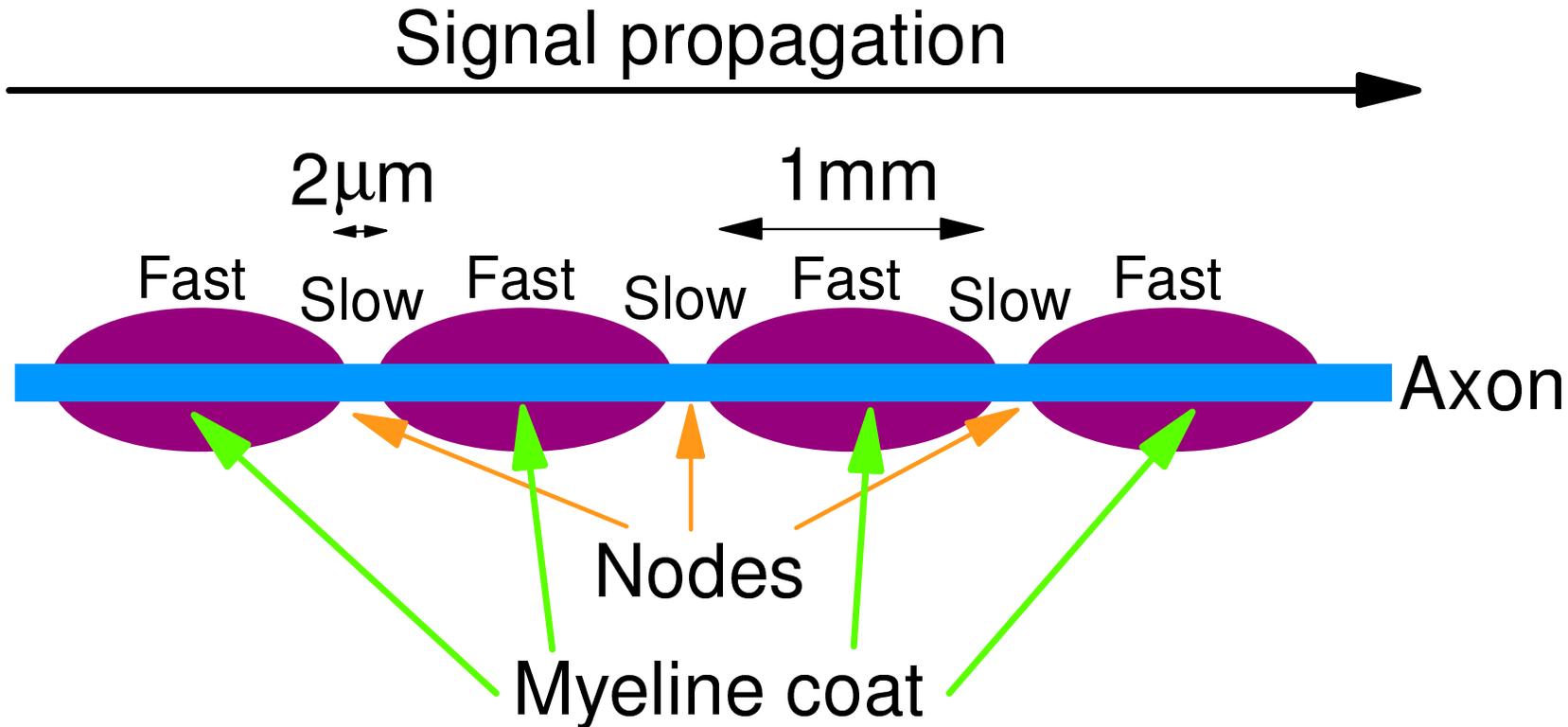
$$\partial_t u(t, x) = \partial_{xx} u(t, x) + f(u(t, x)), \quad x \in \mathbb{R}.$$

$u(x)$

- Many physical models have a discrete spatial structure \rightarrow LDEs.
- No need for α to be large; some models even have $\alpha < 0$.
- Main theme: qualitative differences between PDEs and LDEs.

Signal Propagation through Nerves

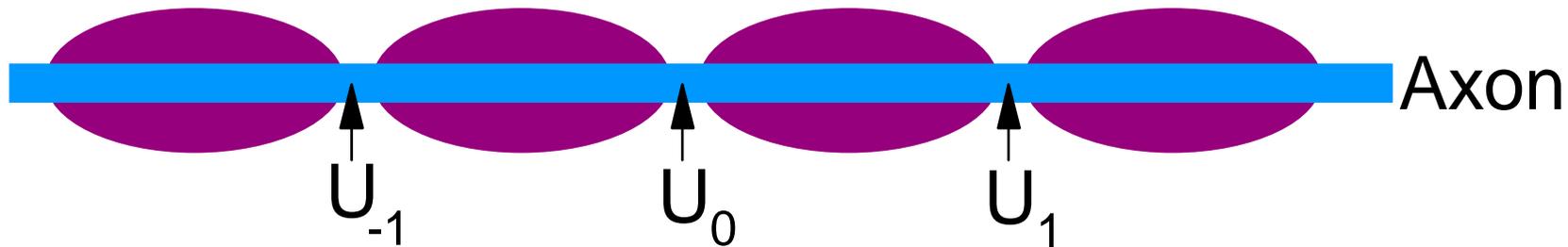
Nerve fibres carry signals over large distances (meter range).



- Fiber has myeline coating with periodic gaps called *nodes of Ranvier* .
- Fast propagation in coated regions, but signal loses strength rapidly (mm-range)
- Slow propagation in gaps, but signal chemically reinforced.

Signal Propagation: The Model

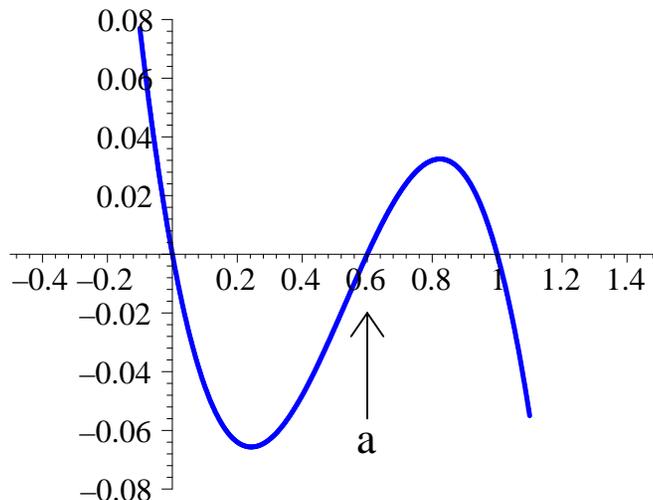
One is interested in the potential U_j at the node sites.



Signals appear to "hop" from one node to the next [Lillie, 1925].

Ignoring recovery, one arrives at the LDE [Keener and Sneyd, 1998]

$$\frac{d}{dt}U_j(t) = U_{j+1}(t) + U_{j-1}(t) - 2U_j(t) + g(U_j(t); a), \quad j \in \mathbb{Z}.$$



Bistable nonlinearity g given by

$$g(u; a) = u(a - u)(u - 1).$$

Signal Propagation: PDE

In continuum limit: Nagumo LDE becomes Nagumo PDE

$$\partial_t u = \partial_{xx} u + u(a - u)(u - 1).$$

Starting step [Fife, McLeod]: travelling waves.

Travelling wave $u(x, t) = \phi(x + ct)$ satisfies:

$$c\phi'(\xi) = \phi''(\xi) + \phi(\xi)(a - \phi(\xi))(\phi(\xi) - 1).$$

Interested in pulse solutions connecting 0 to 1, i.e.

$$\lim_{\xi \rightarrow -\infty} \phi(\xi) = 0, \quad \lim_{\xi \rightarrow +\infty} \phi(\xi) = 1.$$

Signal Propagation: PDE

Recall travelling wave ODE

$$c\phi'(\xi) = \phi''(\xi) + \phi(\xi)(a - \phi(\xi))(\phi(\xi) - 1).$$

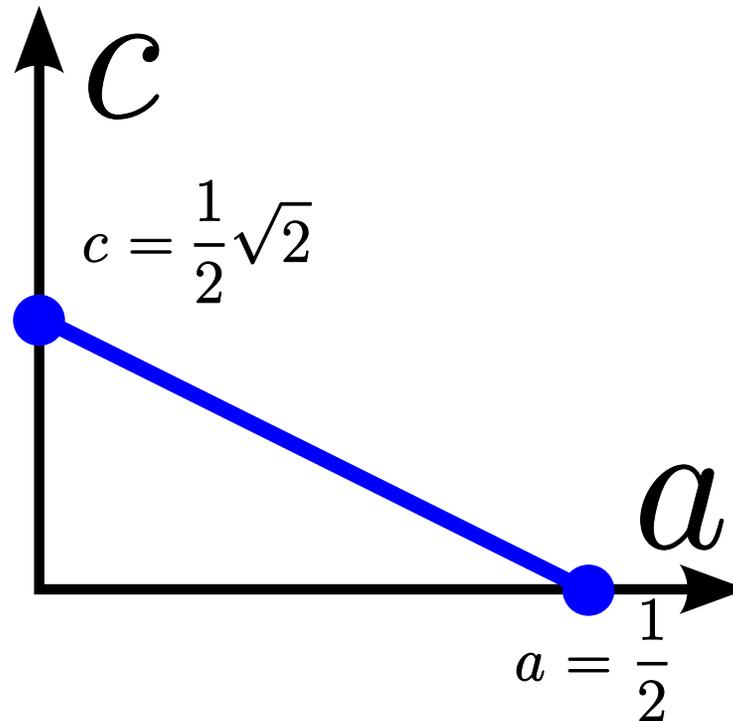
$$\lim_{\xi \rightarrow -\infty} \phi(\xi) = 0,$$

$$\lim_{\xi \rightarrow +\infty} \phi(\xi) = 1.$$

Explicit solutions available:

$$\phi(\xi) = \frac{1}{2} + \frac{1}{2} \tanh\left(\frac{1}{4}\sqrt{2}\xi\right),$$

$$c(a) = \frac{1}{\sqrt{2}}(1 - 2a).$$



Signal Propagation: LDE

Recall the Nagumo LDE

$$\frac{d}{dt}U_j(t) = \frac{1}{h^2}[U_{j+1}(t) + U_{j-1}(t) - 2U_j(t)] + g(U_j(t); a), \quad j \in \mathbb{Z}.$$

Travelling wave profile $U_j(t) = \phi(j + ct)$ must satisfy:

$$c\phi'(\xi) = \frac{1}{h^2}[\phi(\xi + 1) + \phi(\xi - 1) - 2\phi(\xi)] + g(\phi(\xi); a)$$

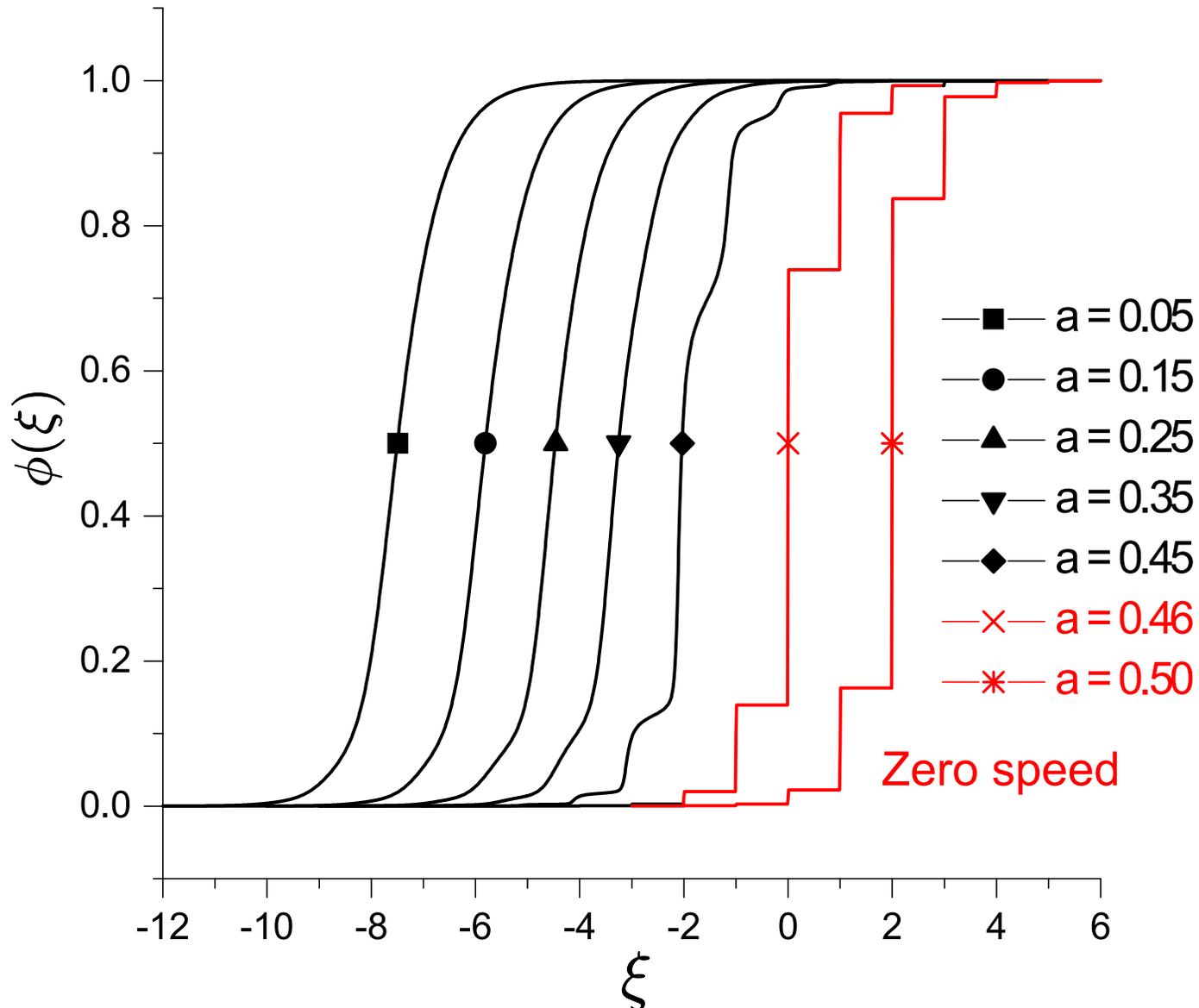
$$\lim_{\xi \rightarrow -\infty} \phi(\xi) = 0,$$

$$\lim_{\xi \rightarrow +\infty} \phi(\xi) = 1.$$

- Notice that wave speed c enters in singular fashion.
- When $c \neq 0$, this is a functional differential equation of mixed type (MFDE).
- When $c = 0$, this is a difference equation.

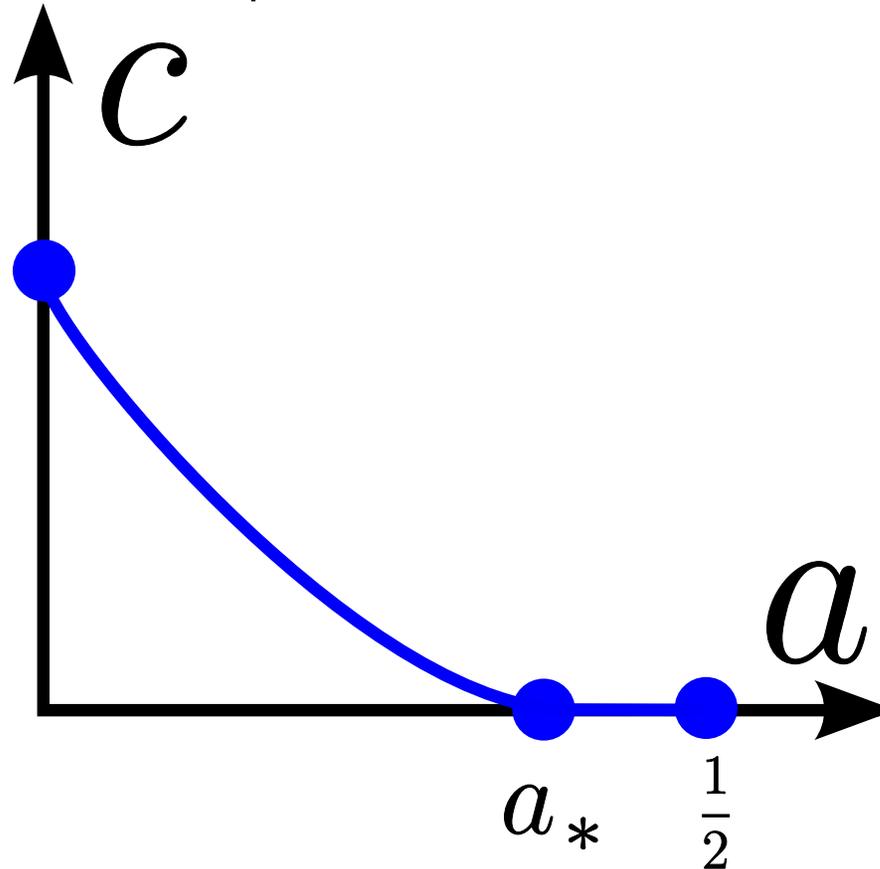
Discrete Nagumo LDE - Propagation failure

Travelling waves for the discrete Nagumo LDE connecting $0 \rightarrow 1$.



Propagation

Typical wave speed c versus a plot for discrete reaction-diffusion systems:



In principle, can have $a_* = \frac{1}{2}$ or $a_* < \frac{1}{2}$.

In case $a_* < \frac{1}{2}$, then we say that LDE suffers from **propagation failure**.

Propagation failure widely studied; pioneered by [Keener].

Signal Propagation: Comparison

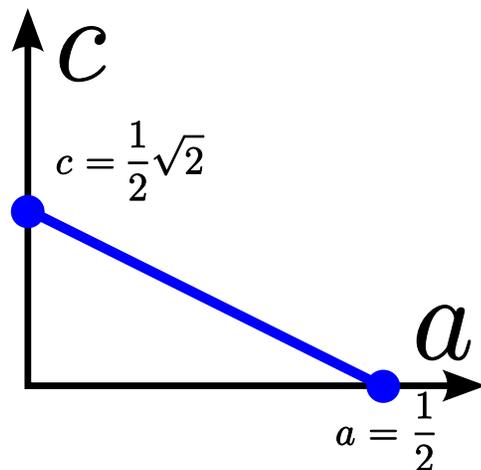
PDE

$$\partial_t u = \partial_{xx} u + g(u, a)$$

Travelling wave $u = \phi(x + ct)$ satisfies:

$$c\phi'(\xi) = \phi''(\xi) + g(\phi(\xi); a)$$

Travelling waves connecting 0 to 1:



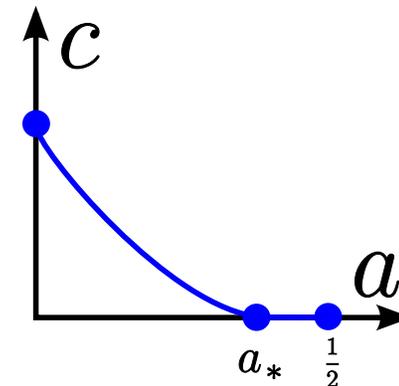
LDE

$$\frac{d}{dt} U_j = U_{j+1} + U_{j-1} - 2U_j + g(U_j; a)$$

Travelling wave $U_j = \phi(j + ct)$ satisfies:

$$c\phi'(\xi) = \phi(\xi + 1) + \phi(\xi - 1) - 2\phi(\xi) + g(\phi(\xi); a)$$

Travelling waves connecting 0 to 1:



Propagation failure if $a_* < \frac{1}{2}$.

Propagation failure

Consider travelling wave MFDE with **saw-tooth** nonlinearity

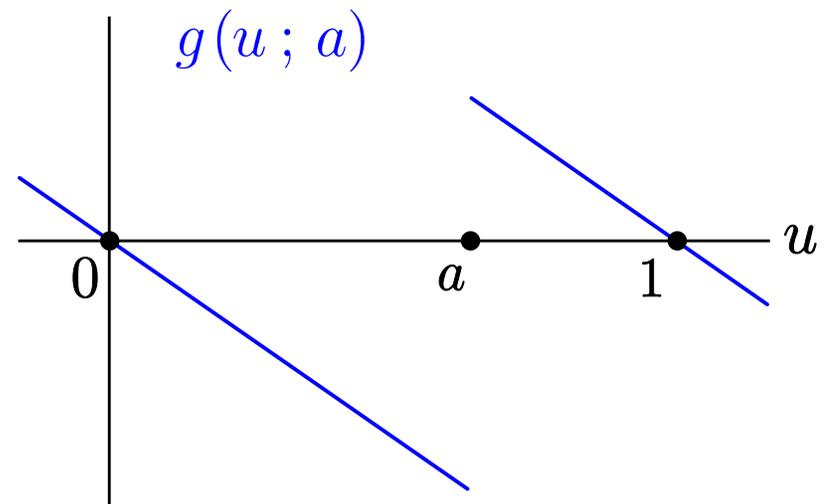
$$c\phi'(\xi) = \frac{1}{h^2}[\phi(\xi + 1) + \phi(\xi - 1) - 2\phi(\xi)] + g(\phi(\xi); a)$$

$$\lim_{\xi \rightarrow -\infty} \phi(\xi) = 0,$$

$$\lim_{\xi \rightarrow +\infty} \phi(\xi) = 1.$$

Thm. [Cahn, Mallet-Paret, Van Vleck]:
Propagation failure for all $h > 0$ (1999).

Linear analysis with Fourier series.



Propagation failure

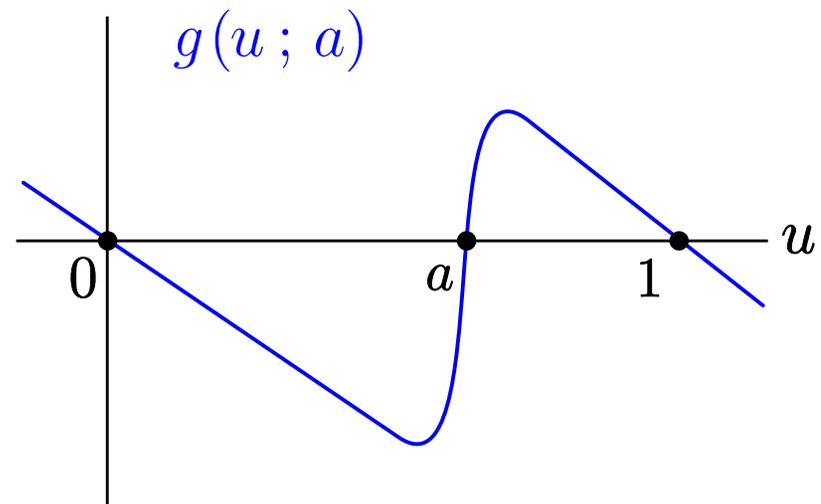
Consider travelling wave MFDE with **near-saw-tooth** nonlinearity

$$c\phi'(\xi) = \frac{1}{h^2}[\phi(\xi + 1) + \phi(\xi - 1) - 2\phi(\xi)] + g(\phi(\xi); a)$$

$$\lim_{\xi \rightarrow -\infty} \phi(\xi) = 0,$$

$$\lim_{\xi \rightarrow +\infty} \phi(\xi) = 1.$$

Thm. [Mallet-Paret]: Propagation failure when g sufficiently close to saw-tooth.



Propagation failure

Consider travelling wave MFDE with **generic** bistable nonlinearity

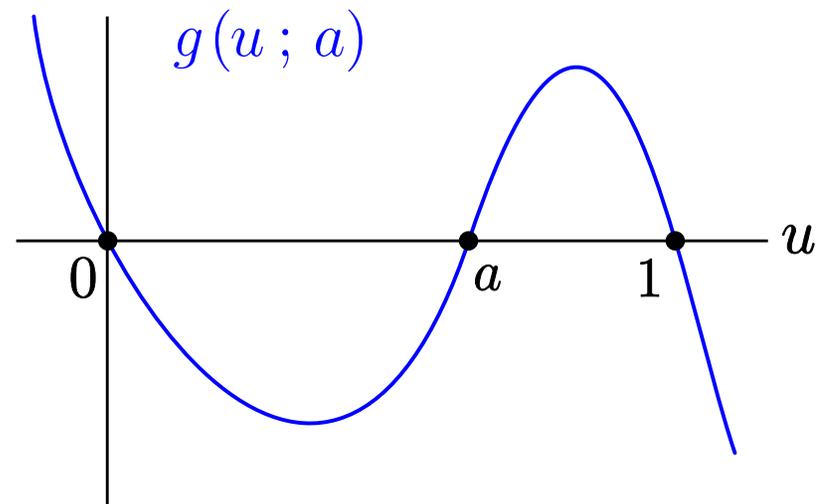
$$c\phi'(\xi) = \frac{1}{h^2}[\phi(\xi + 1) + \phi(\xi - 1) - 2\phi(\xi)] + g(\phi(\xi); a)$$

$$\lim_{\xi \rightarrow -\infty} \phi(\xi) = 0,$$

$$\lim_{\xi \rightarrow +\infty} \phi(\xi) = 1.$$

Thm. [Hoffman, Mallet-Paret]: Generic condition on g guarantees propagation failure.

Unknown if cubic satisfies this condition for all $h > 0$.



Propagation failure

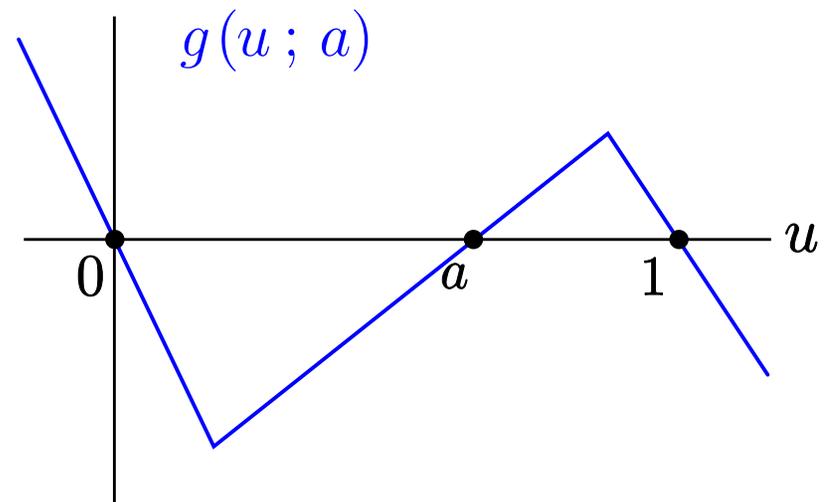
Consider travelling wave MFDE with **zig-zag** bistable nonlinearity

$$c\phi'(\xi) = \frac{1}{h^2}[\phi(\xi + 1) + \phi(\xi - 1) - 2\phi(\xi)] + g(\phi(\xi); a)$$

$$\lim_{\xi \rightarrow -\infty} \phi(\xi) = 0,$$

$$\lim_{\xi \rightarrow +\infty} \phi(\xi) = 1.$$

Thm. [Elmer]: There exist countably many h for which there is **no** propagation failure.



Propagation Failure

Recall travelling wave MFDE:

$$c\phi'(\xi) = \frac{1}{h^2}[\phi(\xi + 1) + \phi(\xi - 1) - 2\phi(\xi)] + g_{\text{cub}}(\phi(\xi); a)$$

$$\lim_{\xi \rightarrow -\infty} \phi(\xi) = 0,$$

$$\lim_{\xi \rightarrow +\infty} \phi(\xi) = 1.$$

When $c = 0$, can restrict to $\xi \in \mathbb{Z}$: recurrence relation!

With $p_j = \phi(j)$ and $r_j = \phi(j + 1)$, we find

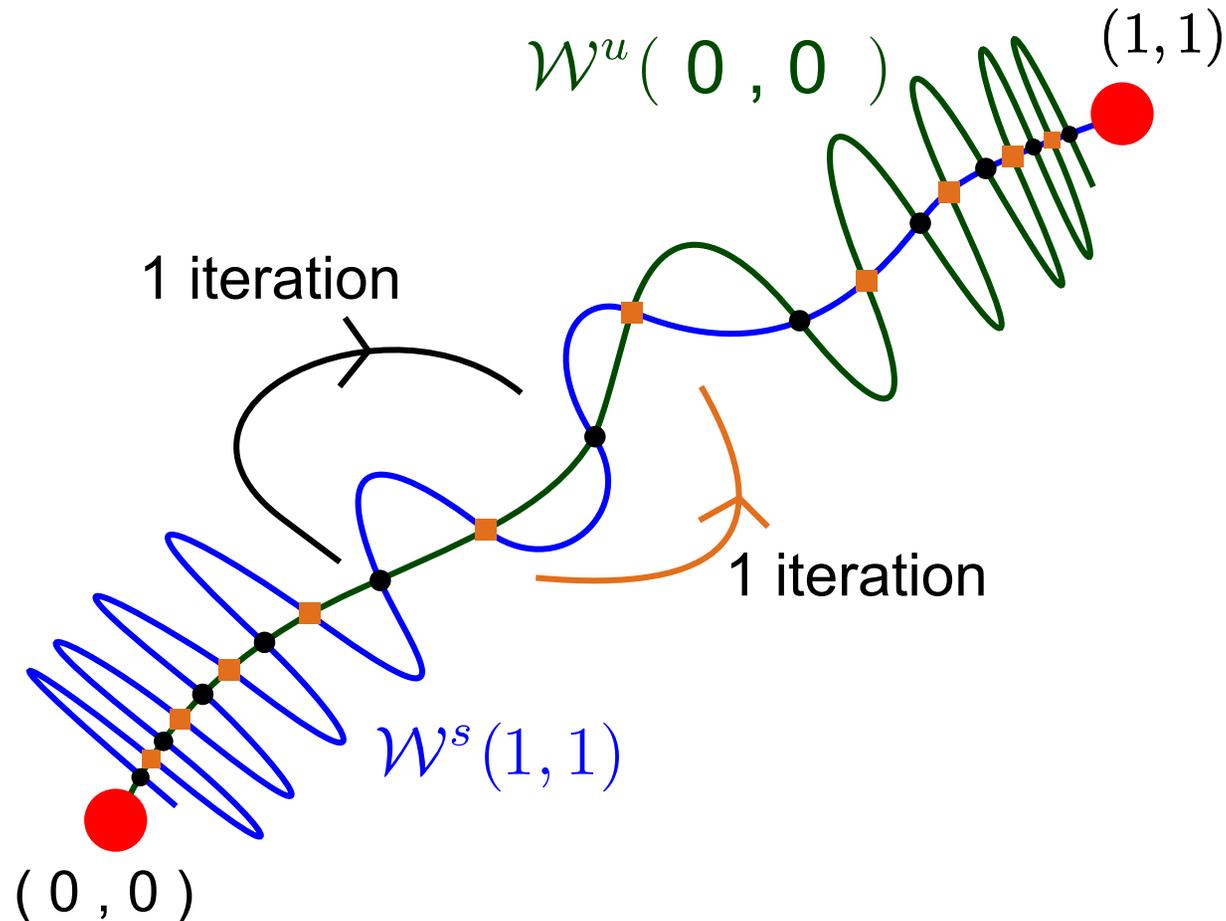
$$\begin{aligned} p_{j+1} &= r_j \\ r_{j+1} &= -p_j + 2r_j - h^2 r_j (r_j - a)(1 - r_j). \end{aligned}$$

Saddles $(0, 0)$ and $(1, 1)$.

Propagation Failure

$$\begin{aligned} p_{j+1} &= r_j \\ r_{j+1} &= -p_j + 2r_j - \alpha^{-1}r_j(r_j - a)(1 - r_j). \end{aligned}$$

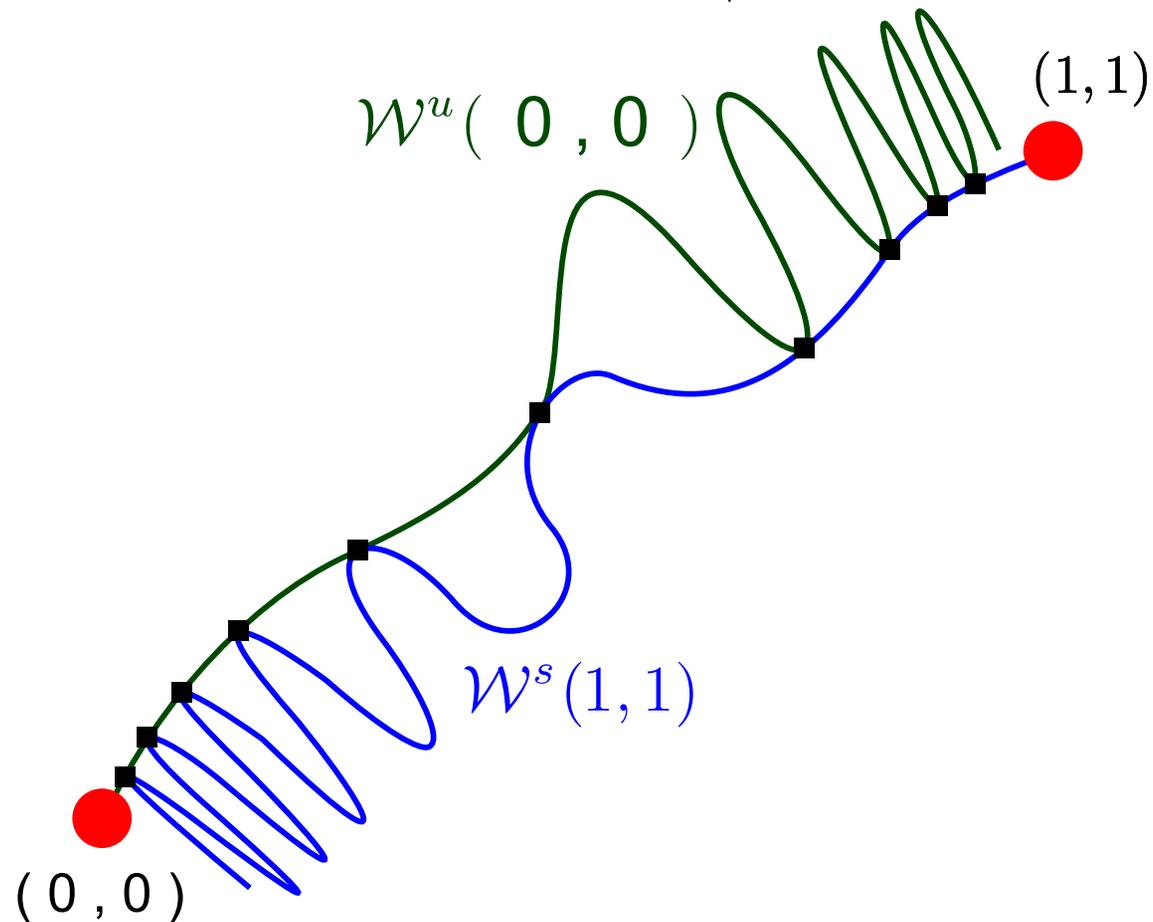
For $a = \frac{1}{2}$, site-centered (orange) and bond-centered (black) solutions. Generically:



Propagation Failure

$$\begin{aligned} p_{j+1} &= r_j \\ r_{j+1} &= -p_j + 2r_j - \alpha^{-1}r_j(r_j - a)(1 - r_j). \end{aligned}$$

Two branches coincide and annihilate at $a = a_*$.



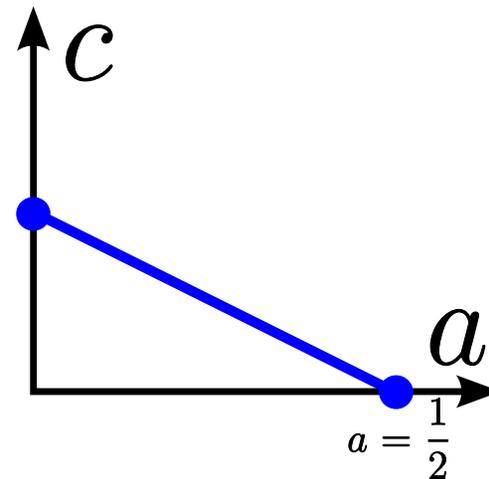
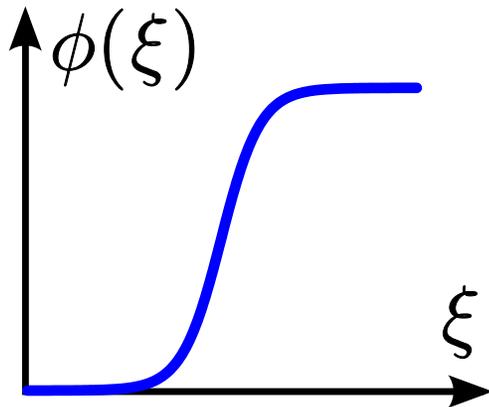
Propagation Failure

Discretizations of cubic may also involve multiple lattice sites:

$$\frac{d}{dt}U_j = \frac{1}{h^2}[U_{j-1} + U_{j+1} - 2U_j] + \frac{1}{2}U_j(U_{j+1} + U_{j-1} - 2a)(1 - U_j).$$

Explicit solutions available:

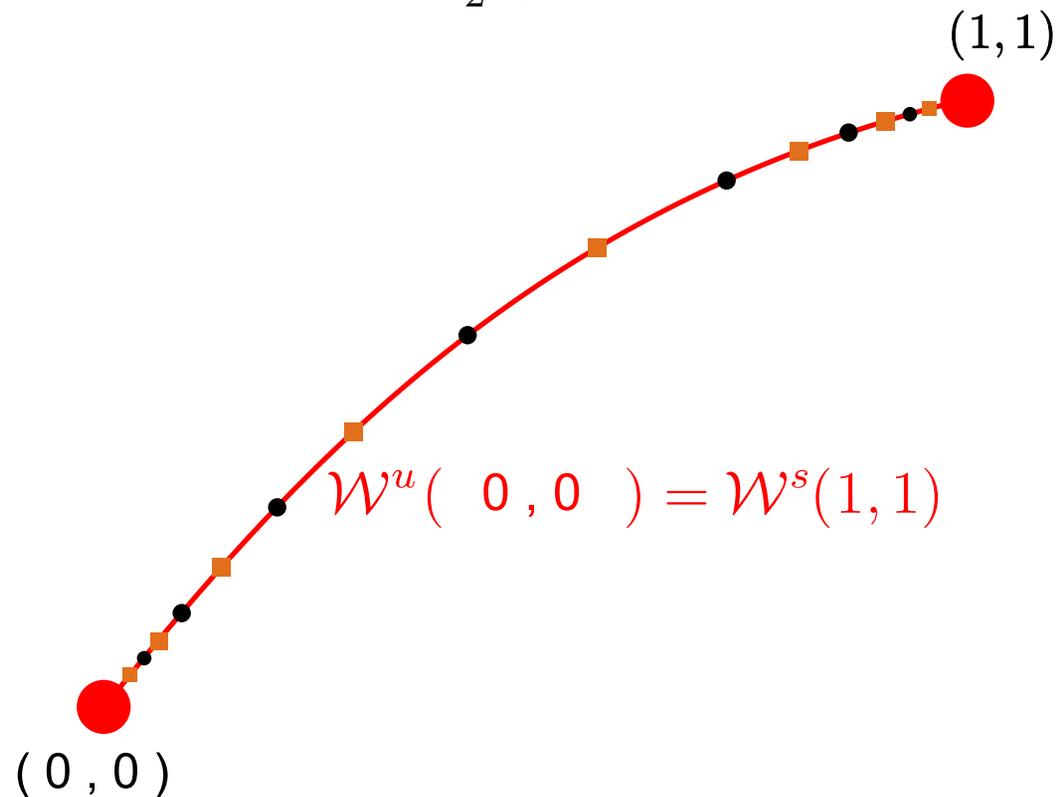
$$U_j(t) = \frac{1}{2} + \frac{1}{2} \tanh \left(\operatorname{arcsinh} \left(\frac{1}{4} \sqrt{2} h \right) (j + ct) \right), \quad c(a) = \frac{(1 - 2a)}{4 \operatorname{arcsinh} \left(\frac{1}{4} \sqrt{2} h \right)}.$$



No propagation failure; smooth wave profile.

Propagation Failure - Discrete map

Smooth standing wave profile at $a = \frac{1}{2}$ gives:



Site centered and bond centered solutions now connected by continuous branch of standing waves.

Q: What happens to manifolds when $a \neq \frac{1}{2}$?

Do intersections disappear (no prop failure) or survive (prop failure)?

Lattice point of view

Let us write LDE as:

$$\frac{d}{dt}U(t) = \mathcal{F}(U(t); a),$$

with $U(t) \in \ell^\infty$ and $\mathcal{F} : \ell^\infty \times [0, 1] \rightarrow \ell^\infty$.

Travelling waves $U_j(t) = \phi(j + ct)$ satisfy some MFDE

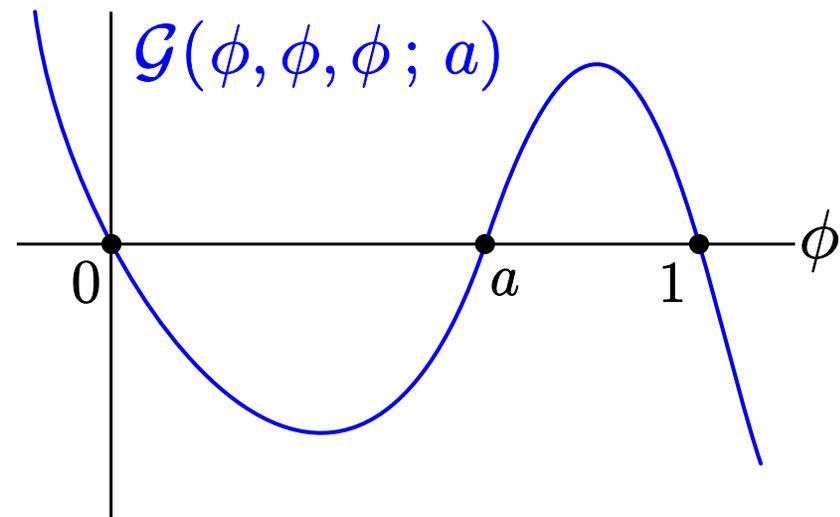
$$c\phi'(\xi) = \mathcal{G}(\phi(\xi - 1), \phi(\xi), \phi(\xi + 1); a).$$

- Assumption: We have

$$\partial_{\phi(\xi-1)}\mathcal{G} > 0,$$

$$\partial_{\phi(\xi+1)}\mathcal{G} > 0.$$

- Assumption: The function $\mathcal{G}(\phi, \phi, \phi; a)$ is bistable. In special case $a = \frac{1}{2}$, it is symmetric.



Lattice point of view

Recall LDE as:

$$\frac{d}{dt}U(t) = \mathcal{F}(U(t); a),$$

and travelling wave MFDE

$$c\phi'(\xi) = \mathcal{G}\left(\phi(\xi - 1), \phi(\xi), \phi(\xi + 1); a\right)$$

Suppose at $a = \frac{1}{2}$ we have a **smooth** solution $p(\xi)$ to

$$0 = \mathcal{G}\left(p(\xi - 1), p(\xi), p(\xi + 1); a\right), \quad \xi \in \mathbb{R}.$$

Then for every $\vartheta \in \mathbb{R}$, we have equilibrium solution $p^{(\vartheta)} \in \ell^\infty$ to our LDE:

$$\mathcal{F}(p^{(\vartheta)}; \frac{1}{2}) = 0, \quad p_j^{(\vartheta)} = p(\vartheta + j).$$

Invariant Manifold

Recall $p^{(\vartheta)} \in \ell^\infty$ with $p_j^{(\vartheta)} = p(\vartheta + j)$.

Notice that

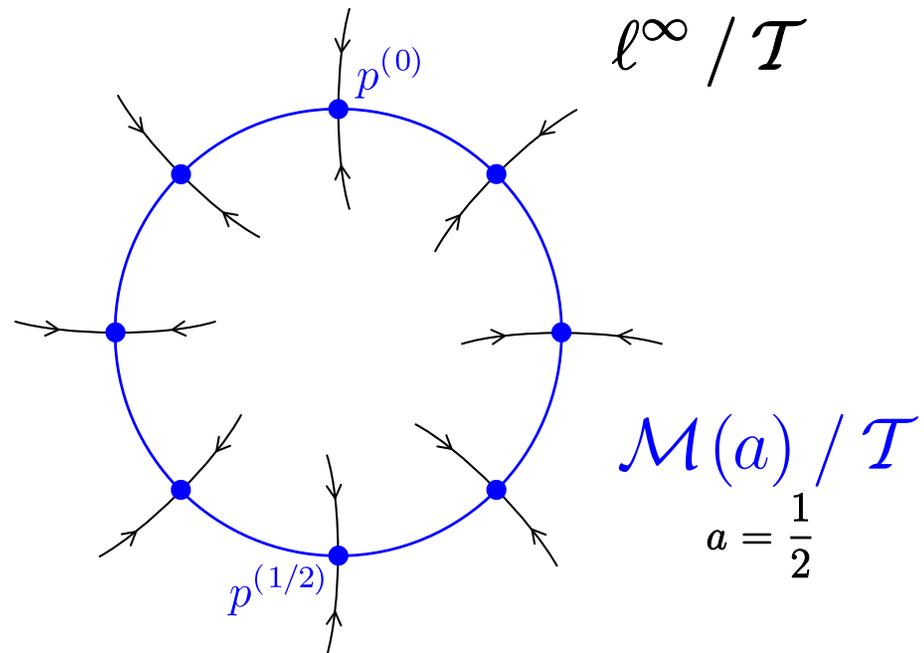
$$p^{(\vartheta)} = \mathcal{T}p^{(\vartheta+1)},$$

where $\mathcal{T} : \ell^\infty \rightarrow \ell^\infty$ is right-shift operator $(\mathcal{T}u)_j = u_{j-1}$.

Combining these equilibria gives a smooth manifold

$$\mathcal{M}(a = \frac{1}{2}) = \{p^{(\vartheta)}\}_{\vartheta \in \mathbb{R}}.$$

After dividing out \mathcal{T} , we get a ring!

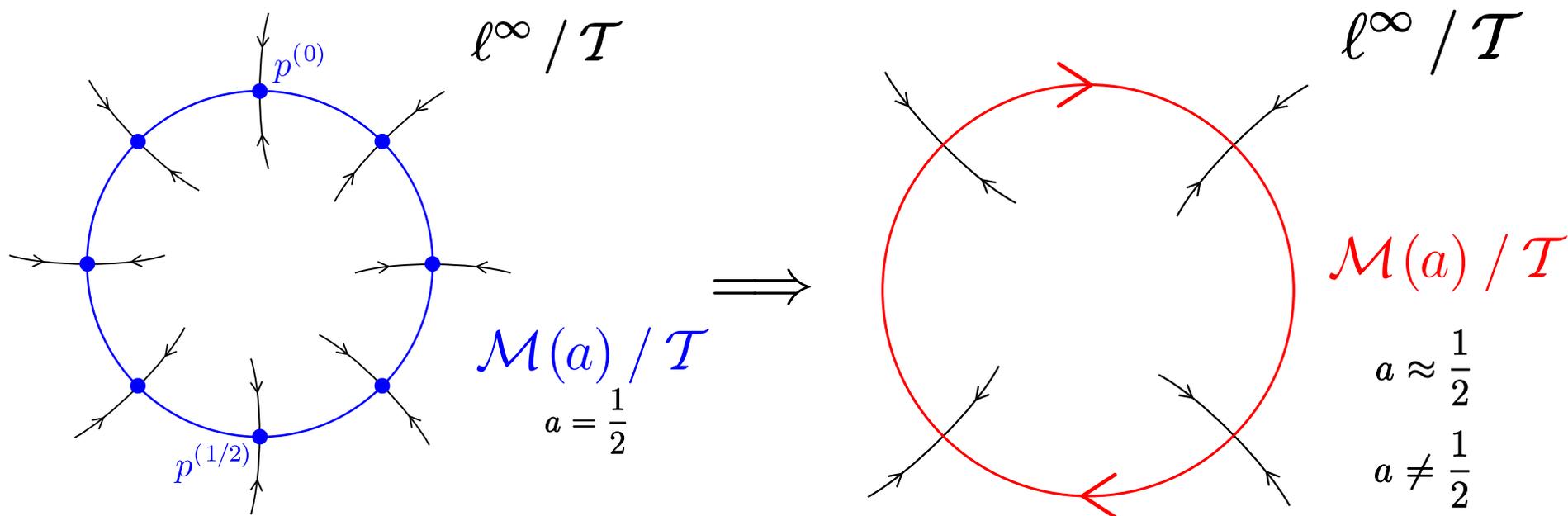


Invariant Manifold - Scenario #1

Based on spectral stability of equilibria $p^{(\vartheta)}$ [Chow, Mallet-Paret, Shen, 1998] and comparison principles can prove:

Prop: The manifold $\mathcal{M}(a = \frac{1}{2})$ is normally hyperbolic.

Possible scenario #1 for persistence of $\mathcal{M}(a)$ with $a \neq \frac{1}{2}$:



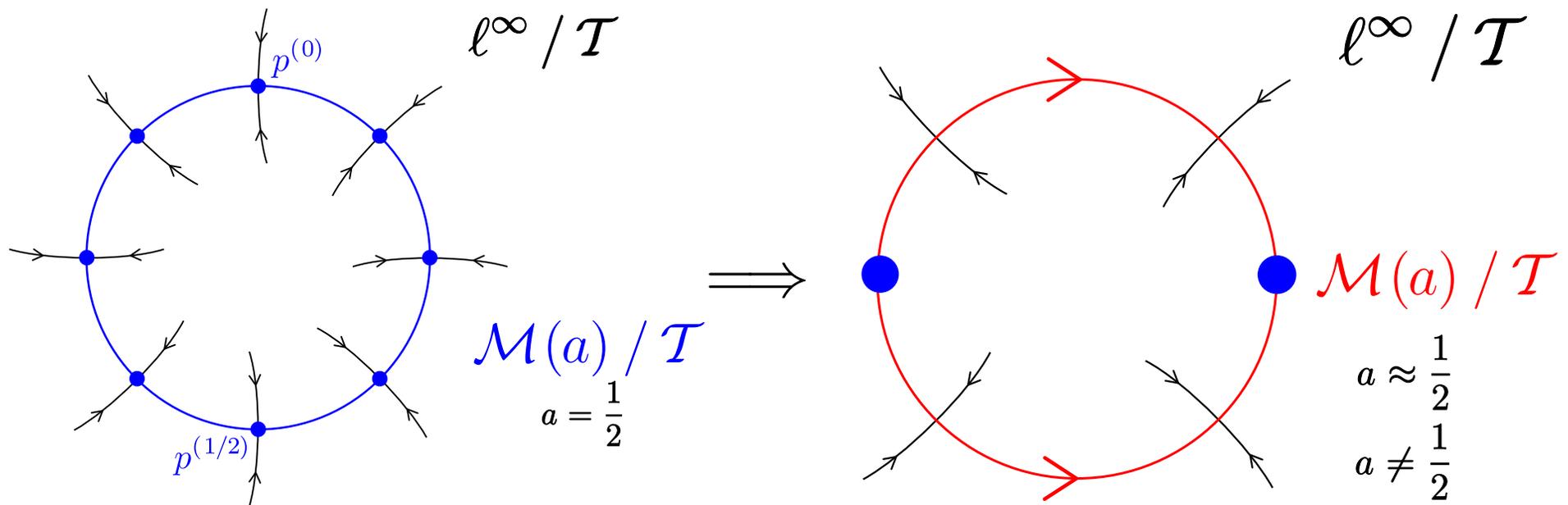
No equilibria survive; $\mathcal{M}(a)$ is orbit of travelling wave. **No Propagation Failure.**

Invariant Manifold - Scenario #2

Based on spectral stability of equilibria $p^{(\vartheta)}$ [Chow, Mallet-Paret, Shen, 1998] and comparison principles can prove:

Prop: The manifold $\mathcal{M}(a = \frac{1}{2})$ is normally hyperbolic.

Possible scenario #2 for persistence of $\mathcal{M}(a)$ with $a \neq \frac{1}{2}$:



One or more equilibria survive. **Propagation Failure***.

*Certain terms and conditions apply...

Dynamics near \mathcal{M}

Angular coordinate θ measures position along $\mathcal{M}(a)$. Dynamics given by

$$\frac{d}{dt}\theta = (a - \frac{1}{2})\Psi(\theta) + O\left(\left|a - \frac{1}{2}\right|^2\right),$$

in which $\Psi(\theta)$ given by

$$\Psi(\vartheta) = \sum_{j \in \mathbb{Z}} q_j^{(\vartheta)} \partial_a \mathcal{G}\left(p_{j-1}^{(\vartheta)}, p_j^{(\vartheta)}, p_{j+1}^{(\vartheta)}; a = \frac{1}{2}\right).$$

Here $q^{(\vartheta)}$ is adjoint eigenvector; i.e. solves $L^{(\vartheta)*}q^{(\vartheta)} = 0$ with

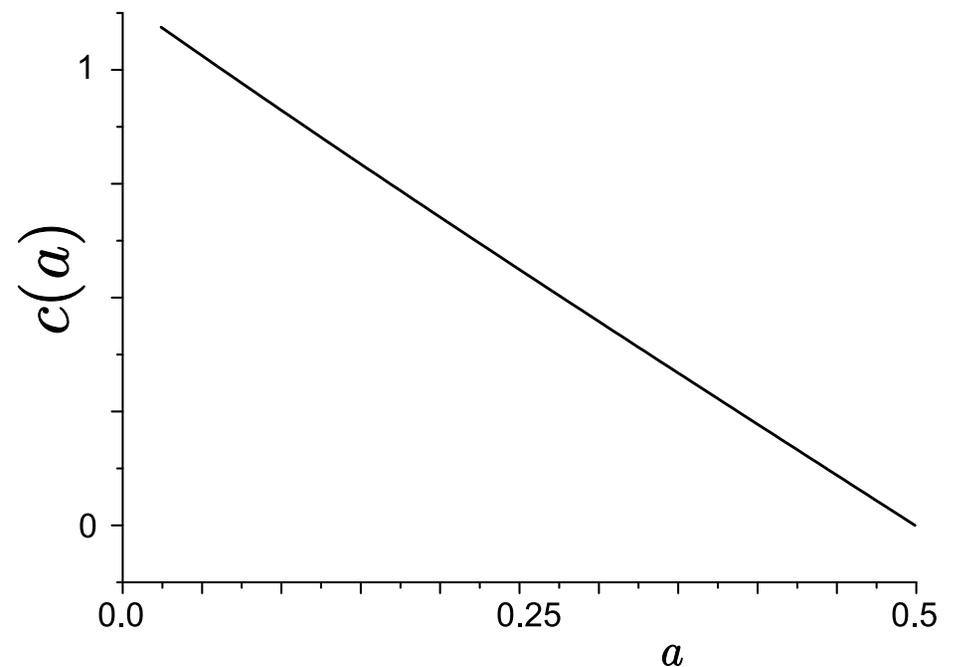
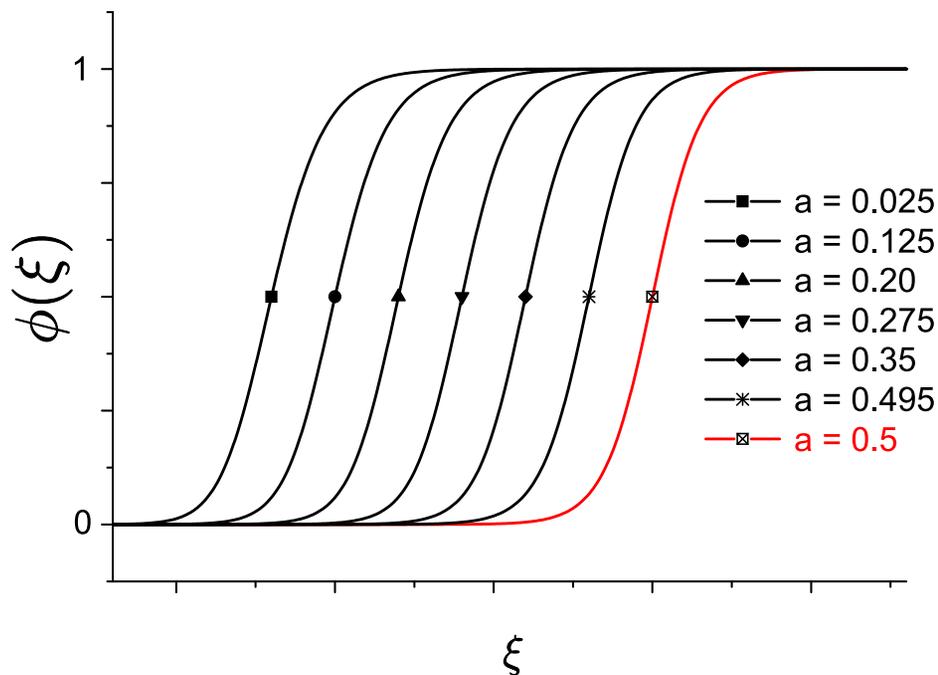
$$\begin{aligned} (L^{(\vartheta)*}w)_j &= \partial_{\phi(\xi-1)} \mathcal{G}\left(p_j^{(\vartheta)}, p_{j+1}^{(\vartheta)}, p_{j+2}^{(\vartheta)}; \frac{1}{2}\right) w_{j+1} \\ &\quad + \partial_{\phi(\xi)} \mathcal{G}\left(p_{j-1}^{(\vartheta)}, p_j^{(\vartheta)}, p_{j+1}^{(\vartheta)}; \frac{1}{2}\right) w_j \\ &\quad + \partial_{\phi(\xi+1)} \mathcal{G}\left(p_{j-2}^{(\vartheta)}, p_{j-1}^{(\vartheta)}, p_j^{(\vartheta)}; \frac{1}{2}\right) w_{j-1}. \end{aligned}$$

Known: $q_j^{(\vartheta)} > 0$ for all $j \in \mathbb{Z}$ and $\vartheta \in \mathbb{R}$. So $\partial_a \mathcal{G} < 0$ guarantees **no** prop failure.

Propagation Failure

Thm. [H., Sandstede, Pelinovsky] **No** prop failure for LDE

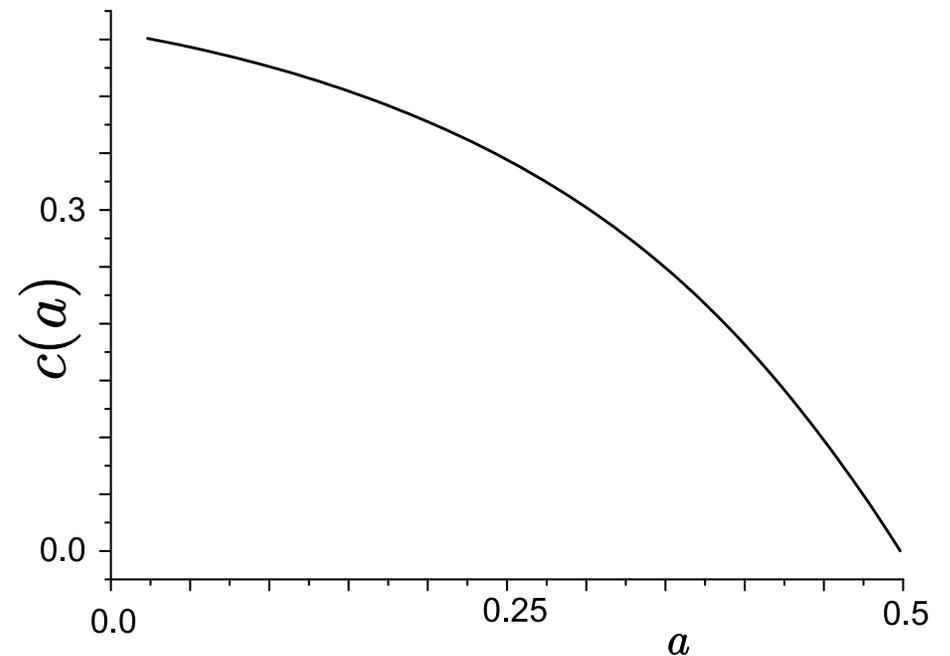
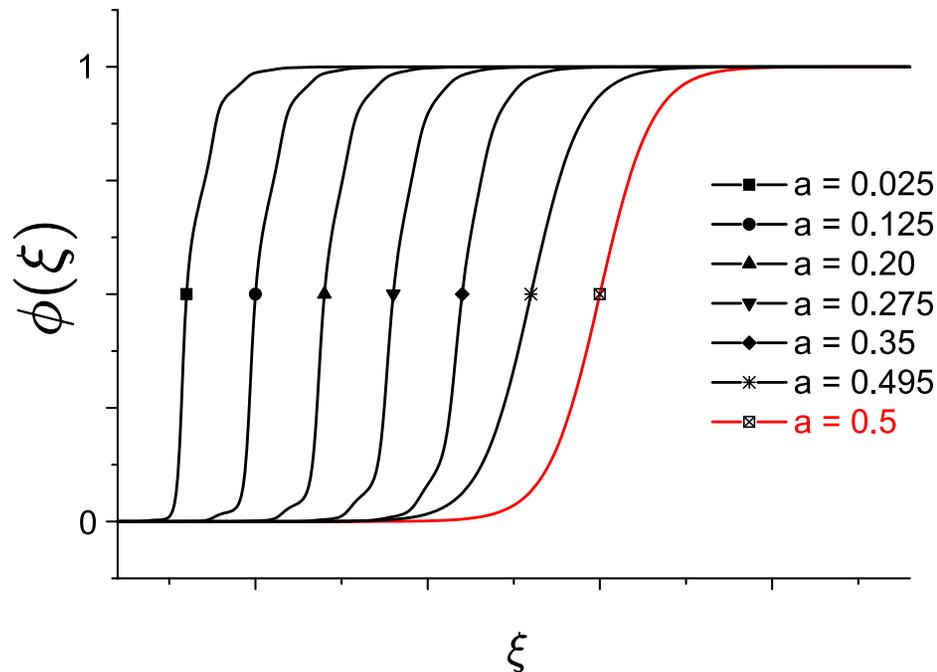
$$\frac{d}{dt}u_j = u_{j-1} + u_{j+1} - 2u_j + (u_j - a) \left(u_{j-1}(1 - u_{j+1}) + u_{j+1}(1 - u_{j-1}) \right)$$



Propagation Failure

Thm. [H., Sandstede, Pelinovsky] **No** prop failure for LDE

$$\frac{d}{dt}u_j = u_{j-1} + u_{j+1} - 2u_j + (u_j - a) \left(u_{j-1}(1 - u_{j+1}) + u_{j+1}(1 - u_{j-1}) \right) - \frac{5}{4} \left(a - \frac{1}{2} \right) \sin(2\pi u_j).$$

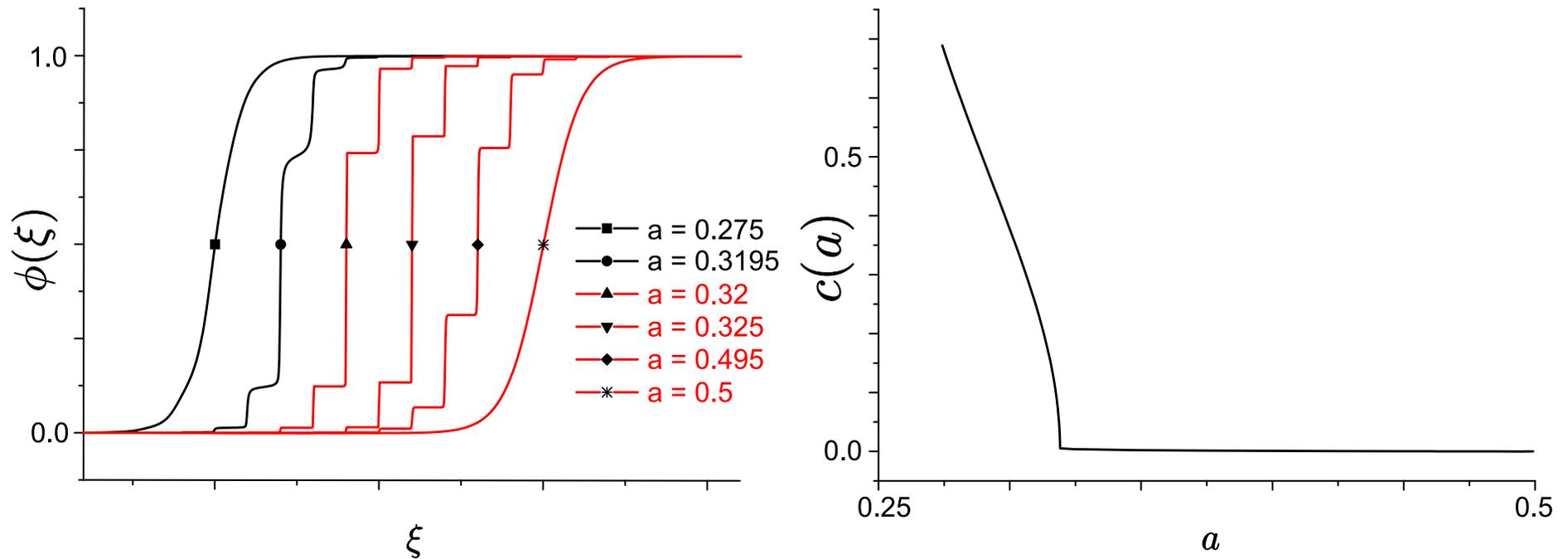


Here $\partial_a \mathcal{G}$ may have both signs, but (numerically) $\Psi(\theta) < 0$ for all θ .

Propagation Failure

Thm. [H., Sandstede, Pelinovsky] **Do** have prop failure for LDE

$$\frac{d}{dt}u_j = u_{j-1} + u_{j+1} - 2u_j + 4u_j(1 - u_j)(u_{j-1} + u_{j+1} - 2a) - 5\left(a - \frac{1}{2}\right) \sin(2\pi u_j) \left(\frac{6}{5} + \frac{8}{5}u\right).$$



Numerically computed: $\Psi(\theta = 0) < 0 < \Psi(\theta = \frac{1}{2})$.

Discussion

Recall PDE $u_t = u_{xx} + g(u; a)$.

- Active interest in multi-lattice-site discretizations of g that admit continuous branch of stationary solutions [Barashenkov, Oxtoby, Pelinovsky, Dmitriev, Kevrekidis, Yoshikawa].
- One generally expects size of propagation failure interval to be exponentially small in h .
- For higher dimensional problems, indications are that using 'small enough' $h > 0$ to reduce influence of propagation failure can hurt numerical performance [Beyn, Speight].